# NEW RESULTS on SQUARED BESSEL PARTICLE SYSTEMS and WISHART PROCESSES

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Random matrices and their applications Kyoto 21-25/5/2018

### Articles

P. Graczyk, J. Małecki On Squared Bessel particle systems, to appear in Bernoulli, 2018

- K. Bogus, P. Graczyk, J. Małecki Properties of β-Squared Bessel particle systems, preprint, 2018
- P. Graczyk, J. Małecki, E. Mayerhofer

A Characterization of Wishart Processes and Wishart Distributions,

Stoch.Proc. Appl., 128 (2018), 1347–1385.

P. Graczyk, J. Małecki,

Strong solutions of non-colliding particle systems. Electron. J. Probab. 19 (2014), pp. 1–21. **Objective:** study solutions of the SDEs system for BESQ particles

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Göing-Jaeschke, Yor [Bernoulli 2003]:  $\alpha \in \mathbb{R}, x \in \mathbb{R}, X \in \mathbb{R}$ 

Relation with Squared Bessel matrix (Wishart) processes

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If  $N_t =$  Brownian Motion on  $p \times \alpha$  matrices  $(\alpha \in \mathbb{N})$ , define  $\boxed{\mathbb{Y}_t = N_t N_t^T \in \overline{S}_p^+, t \ge 0.}$  Then, by Itô formula, for  $\alpha \ge n$ , Recall classical Squared Bessel Matrix(Wishart) processes on  $\bar{S}_p^+$ :  $S_p$  = the space of symmetric  $p \times p$  matrices,  $\bar{S}_p^+$  = the cone of positive semi-definite matrices

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 $d\mathbb{Y}_t = \sqrt{\mathbb{Y}_t} d\mathbb{W}_t + d\mathbb{W}_t^T \sqrt{\mathbb{Y}_t} + \alpha \mathbb{I} dt,$ 

where  $\mathbb{W}_t$  is a Brownian  $p \times p$  matrix(Bru 1991). Bru showed that this matrix SDE has solutions for  $\alpha \in [p-1,\infty) \cup \{1,\ldots,p-1\}$ .

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# Difficulties in solving SDEs for BESQ particles

In equations for non-colliding BESQ particles

 $\begin{aligned} dX_i &= 2\sqrt{|X_i|} dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{X_i - X_j} \mathbb{1}_{\{X_i \neq X_j\}}\right) dt, \quad i = 1, \dots, p \\ X_1(t) &\leq X_2(t) \leq \ldots \leq X_p(t), \qquad t \geq 0, \end{aligned}$ 

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two difficulties appear

- non-Lipschitz functions √x in martingale parts (Yamada-Watanabe th. is 1-dimensional!)
- The drift part contains singularities  $(X_i X_j)^{-1}$ (physicists want to start from (0, ..., 0)!)

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**Proof** (and proofs of other reported results) based on:

- SDEs for symmetric polynomials (explained now)
- "Glueing" solutions for subsystems (explained later on examples)

The elementary symmetric polynomials of  $X = (X_1, ..., X_p)$  are defined, for n = 1, ..., p by  $e_n(X) = \sum_{i_1 < i_2 < ... < i_n} X_{i_1} X_{i_2} \cdot ... \cdot X_{i_n}$ 

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Piotr Graczyk New results on BESQ particle systems and Wishart processes

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The SDEs for symmetric polynomials are non-singular!!!

Methodology introduced for general particle systems in:

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In this research we use extensively main results and methods of EJP(2014).

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– eigenvalue processes of matrix processes on  $\mathit{Sym}_{p\times p}$  verifying the matrix SDE

$$dX_t = h(X_t)dW_tg(X_t) + g(X_t)dW_t^Th(X_t) + b(X_t)dt$$

where the functions  $g, h, b : \mathbb{R} \to \mathbb{R}$  act spectrally on  $Sym_{p \times p}$ .

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we prove, when starting from  $\lambda_1(0) \leq \ldots \leq \lambda_p(0)$ and under natural conditions on the coefficients  $\sigma_i$ ,  $H_{ij}$ ,  $b_i$ (symmetry of H, regularity, non-explosion, non-collision conditions)

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- non-colliding of solutions of this system
- by methods of classical Itô calculus

#### Using polynomials to prove existence for particles SDEs

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- The symmetry of  $H_{ij}(x, y) \Rightarrow$  the singularities  $\frac{1}{\lambda_i \lambda_i}$  cancel!
- Analogous phenomenon occurs for other basic symmetric polynomials of (λ<sub>1</sub>,..., λ<sub>p</sub>)

# Example of SDEs for symmetric polynomials: BESQ particle systems, p = 4

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$$de_{3} = 2\sqrt{e_{3}e_{2} + 6e_{1}e_{4}}dU_{3} + 2(\alpha - 2)e_{2}dt$$
  

$$de_{4} = 2\sqrt{e_{4}e_{3}}dU_{4} + (\alpha - 3)e_{3}dt$$

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(Brownian motions  $U_i$  are not independent, but their brackets  $\langle U_i, U_j \rangle$  are determined)

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   ((-1)<sup>k</sup>e<sub>k</sub>(Λ) is the coefficient of x<sup>p-k</sup> in
   P(x) = Π<sup>p</sup><sub>i=1</sub>(x − λ<sub>i</sub>))

- We solve SDEs for polynomials
- Fact. The map  $e = (e_1, \ldots, e_p) : C_+ \to \mathbb{R}^p$  is a diffeomorphism and extends continuously to  $\overline{C_+}$  ( $(-1)^k e_k(\Lambda)$  is the coefficient of  $x^{p-k}$  in  $P(x) = \prod_{i=1}^p (x \lambda_i)$ )
- we define  $\lambda = \lambda(e_1, \dots, e_p)$  and show that they are solutions of the SDEs system for  $\Lambda$

Application of EJP(2014) to BESQ particles  $dX_i = 2\sqrt{|X_i|}dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{X_i - X_j} \mathbb{1}_{\{X_i \neq X_j\}}\right)dt$ 

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Let  $|\alpha| \in \mathbb{R}^+ \setminus \{0, 1, \dots, p-2\}$ . Then the BESQ particle system has a unique non-colliding solution for t > 0. If  $\alpha \ge p-1$  and  $X_1(0) \ge 0$ , then the solution is non-negative, i.e.  $X_1(t) \ge 0$ . Application of EJP(2014) to BESQ particles  $dX_i = 2\sqrt{|X_i|}dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{X_i - X_j} \mathbb{1}_{\{X_i \neq X_j\}}\right)dt$ 

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One of assumptions of EJP(2014) fails if  $|\alpha| \in \{0, 1, \dots, p-2\}$ 

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But strong non-colliding solution exists!

Thanks to the results of EJP(2014), it is enough to prove the existence of a non-colliding weak solution. We prove it by glueing non-negative and non-positive solutions. Let us explain it on the simplest example. Existence of non-colliding BESQ particle systems for  $\alpha \in \{0, 1, ..., p - 2\}$   $dX_i = 2\sqrt{|X_i|}dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{X_i - X_j} \mathbb{1}_{\{X_i \neq X_j\}}\right)dt$ Example:  $p = 2, \alpha = 0$ , start from (0, 0)

The system is

$$dX_1 = 2\sqrt{|X_1|}dB_1 + \frac{|X_1| + |X_2|}{X_1 - X_2} \mathbb{1}_{\{X_1 \neq X_2\}}dt,$$
  
$$dX_2 = 2\sqrt{|X_2|}dB_2 + \frac{|X_1| + |X_2|}{X_2 - X_1} \mathbb{1}_{\{X_1 \neq X_2\}}dt,$$

where  $X_1(t) \le X_2(t)$ , t > 0 and  $X(0) = (x_1, x_2)$ .

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$$dX_{2} = 2\sqrt{|X_{2}|}dB_{2} + \frac{|X_{1}| + |X_{2}|}{X_{2} - X_{1}} \mathbb{1}_{\{X_{1} \neq X_{2}\}}dt,$$

where  $X_1(t) \le X_2(t)$ , t > 0 and  $X(0) = (x_1, x_2)$ .

The process  $X_1(t) = X_2(t) \equiv 0$  is a strong solution of the system and it is colliding.

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The processes  $\tilde{X}_1$ ,  $\tilde{X}_2$  are two independent squared Bessel processes of dimension -1 and +1 respectively.

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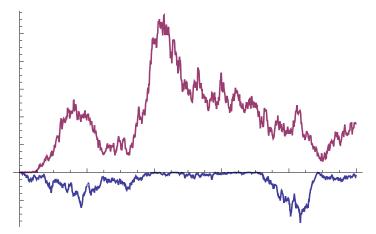
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By independence, these processes do not collide after the start. This illustrates our method of constructing of a non-colliding strong solution for  $\alpha \in \{0, ..., p-2\}$ , by glueing solutions in lower dimensions.

## Example: p = 2, $\alpha = 0$ , start from (0, 0)



Unique non-colliding solution starting from x = (0,0). (Another (non-negative) solution:  $X_1 = X_2 = 0$ .)

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Let rank<sup>+</sup>(x) =  $\sum_{i=1}^{p} 1_{(0,\infty)}(x_i)$ , rank<sup>-</sup>(x) =  $\sum_{i=1}^{p} 1_{(-\infty,0)}(x_i)$ 

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#### Theorem

Pathwise uniqueness with the initial condition X(0) = x, where  $x = (x_1, ..., x_p)$ , holds if and only if one of the following conditions holds

(a) 
$$|\alpha| \notin \{0, 1, \dots, p-2\}$$
  
(b)  $|\alpha| \in \{0, 1, \dots, p-2\}$  and  $(rank^+(x) \ge \frac{p+\alpha-1}{2} \text{ or } rank^-(x) \ge \frac{p-\alpha-1}{2}).$ 

Then there exists unique strong solution, which is non-colliding.

Non–uniqueness for BESQ particle systems  $dX_i = 2\sqrt{|X_i|}dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{X_i - X_j} \mathbb{1}_{\{X_i \neq X_j\}}\right)dt$ 

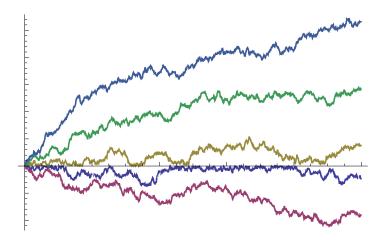
When  $|\alpha| \in \{0, 1, ..., p-2\}$ , rank<sup>+</sup>(x) <  $(p + \alpha - 1)/2$  and rank<sup>-</sup>(x) <  $(p - \alpha - 1)/2$ , then the uniqueness of the strong solutions is violated, i.e. there exist at least two solutions.

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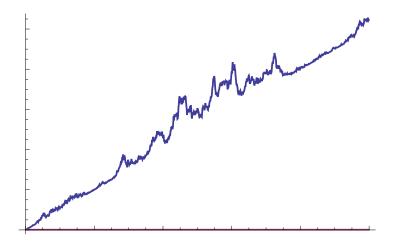
The solutions are all colliding, except one.

 $dX_i = 2\sqrt{|X_i|}dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{|X_i - X_j|} \mathbf{1}_{\{X_i \neq X_j\}}\right) dt$  $p = 5, \ \alpha = 1 \text{ start from } x = (0, \dots, 0)$ 



Unique non-colliding solution  $(BESQ_{nc}^{(-2)}(0,0), BESQ_{nc}^{(3)}(0,0,0))$ 

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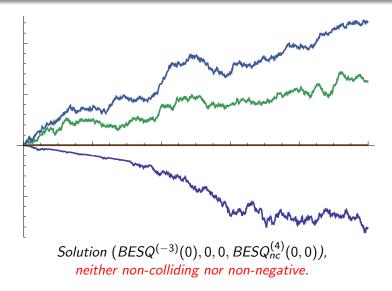


## Unique non-negative solution $(0, 0, 0, 0, BESQ^{(5)}(0))$

Piotr Graczyk

New results on BESQ particle systems and Wishart processes

 $dX_i = 2\sqrt{|X_i|}dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{|X_i - X_j|} \mathbf{1}_{\{X_i \neq X_j\}}\right)dt$  $p = 5, \ \alpha = 1 \text{ start from } x = (0, \dots, 0)$ 



Existence and uniqueness of non-negative solutions  $dX_{i} = 2\sqrt{|X_{i}|}dB_{i} + \left(\alpha + \sum_{j\neq i} \frac{|X_{i}| + |X_{j}|}{X_{i} - X_{j}} \mathbb{1}_{\{X_{i} \neq X_{j}\}}\right)dt$   $0 \leq X_{1}(t) \leq \ldots \leq X_{p}(t)$ 

#### Theorem

There exists a **non-negative** solution with the initial condition X(0) = x, where  $x = (x_1, ..., x_p)$  and  $x_1 \ge 0$ , if and only if one of the following conditions holds

(a) 
$$\alpha \ge p-1$$

(b) 
$$\alpha \in \{0, 1, ..., p-2\}$$
 and  $rk(x) \le \alpha$ .

Then pathwise uniqueness among non-negative solutions holds and there exists unique non-negative strong solution. Existence and uniqueness of non-negative solutions  $dX_{i} = 2\sqrt{|X_{i}|}dB_{i} + \left(\alpha + \sum_{j \neq i} \frac{|X_{i}| + |X_{j}|}{X_{i} - X_{j}} \mathbb{1}_{\{X_{i} \neq X_{j}\}}\right)dt$   $0 \leq X_{1}(t) \leq \ldots \leq X_{p}(t)$ 

#### Theorem

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The proof uses drifts of symmetric polynomials and is the same as for an analogous **non-negativity problem for the BESQ matrix** (Wishart) process: **Stochastic Gindikin set** 

Consider  $\mathbb{Y}_t$ , solution of  $d\mathbb{Y}_t = \sqrt{|\mathbb{Y}_t|} d\mathbb{W}_t + d\mathbb{W}_t^T \sqrt{|\mathbb{Y}_t|} + \alpha \mathbb{I} dt, \quad \alpha \in \mathbb{R}$ 

Problem. For which  $\alpha$  and  $\mathbb{Y}_0$  does the process  $\mathbb{Y}_t$  stay in  $\mathcal{S}_{\rho}^+$ ? (this is the Stochastic Gindikin Set)

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Problem. For which  $\alpha$  and  $\mathbb{Y}_0$  does the process  $\mathbb{Y}_t$  stay in  $\mathcal{S}_p^+$ ? (this is the Stochastic Gindikin Set) Solution in: P. Graczyk, J. Małecki, E. Mayerhofer *A Characterization of Wishart Processes and Wishart Distributions* Stoch.Proc. Appl., 2017

Answer:  $\alpha \ge p-1$  or  $\alpha \in \{0, 1, \dots, p-2\}$  and  $\operatorname{rank}(\mathbb{Y}_0) \le \alpha$ 

Recall that the symmetric polynomials  $e_n = e_n(X)$ , n = 1, ..., p are semimartingales satisfying the following system of SDEs

$$\begin{aligned} de_1 &= 2\sqrt{e_1}dV_1 + p\alpha dt, \\ de_n &= M_n(e_1, \dots, e_p)dV_n + (p - n + 1)(\alpha - n + 1)e_{n-1}dt, \\ de_p &= 2\sqrt{e_{p-1}e_p}dV_p + (\alpha - p + 1)e_{p-1}dt, \end{aligned}$$

where  $V_n$ , n = 1, ..., p are one-dimensional Brownian motions and the functions  $M_n$  are continuous on  $\mathbb{R}^p$ .

# Proof of the Stochastic Gindikin set for $\alpha \in \{0, 1, \dots, p-2\}$

Suppose that a solution  $X_i(t) \ge 0$ . Consider  $\mathbb{E}e_n(t)$  where  $n = \alpha + 1$ . Then

$$\mathbb{E}e_n(t) = e_n(0) + (p - n + 1)(\alpha - n + 1) \int_0^t \mathbb{E}e_{n-1}(s)ds = e_n(0)$$
  
$$\mathbb{E}e_{n+1}(t) = e_{n+1}(0) + (p - n)(\alpha - n)e_n(0)t$$

If  $e_n(0) > 0$ , then the leading term of  $\mathbb{E}e_{n+1}(t)$  is negative and thus  $\mathbb{E}e_{n+1}(t) < 0$  for large t. It implies  $e_n(0) = 0$ , i.e.  $rank(x_0) \le n - 1 = \alpha$ .

#### Application of Stochastic Gindikin Set: First proof of the

### characterization of the Non-Central Gindikin Set (NCGS)

Application of Stochastic Gindikin Set: First proof of the characterization of the Non-Central Gindikin Set (NCGS)

i.e. the parameter set of non-central Wishart Distributions  $\Gamma_{p,\Sigma}(\beta, \mathbf{x})$  on  $\overline{\mathcal{S}_p^+}$  defined by their Laplace transform at  $u \in \mathcal{S}_p^+$  ( $\Sigma \in \mathcal{S}_p^+$  fixed )

$$\mathcal{L}(u) = (\det(I + \Sigma u))^{-\beta} \exp[-\mathrm{Tr}(u(I + \Sigma u)^{-1}x)],$$

Application of Stochastic Gindikin Set: First proof of the **characterization of the Non-Central Gindikin Set (NCGS)** i.e. the parameter set of non-central Wishart Distributions  $\Gamma_{p,\Sigma}(\beta, x)$  on  $\overline{\mathcal{S}_p^+}$  defined by their Laplace transform at  $u \in \mathcal{S}_p^+$ ( $\Sigma \in \mathcal{S}_p^+$  fixed )

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- Attempt by harmonic analysis: Letac,Massam (J. Multiv.Anal. 2008, gap)

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(this is an extrapolation of the Laplace transform of  $Y = \xi_1 \xi_1^T + \ldots + \xi_n \xi_n^T = q(\xi)$  where  $\xi_i \sim N(m_i, \Sigma/2)$  are independent normal vectors in  $\mathbb{R}^p$  and  $x = q(m_1, \ldots, m_n), \beta = \frac{n}{2}$ )

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- Attempt by harmonic analysis: Letac,Massam (J. Multiv.Anal. 2008, gap)

We show that

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BESQ matrix processes are affine processes: the exponent of the Laplace transform of  $\mathbb{Y}_t$  is affine function of the starting point  $x_0$ 

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## THANK YOU FOR YOUR ATTENTION

#### **ARIGATO GOZAIMASU!**

Piotr Graczyk New results on BESQ particle systems and Wishart processes

Multi-indexed BESQ Matrix (Wishart) processes exist for positive integer indices  $\underline{\alpha} = (n_1, \ldots, n_p)$ , and are sums of independent classical Wishart processes.

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Q1. How to define Multi-indexed BESQ Matrix (Wishart) processes for any admissible multi-index  $\underline{\alpha}$ ?

Q2. Study the corresponding multi-indexed BESQ particle systems

Methodology of proofs: SDEs for symmetric polynomials of  $dX_i = 2\sqrt{|X_i|}dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{X_i - X_j} \mathbb{1}_{\{X_i \neq X_j\}}\right)dt$ 

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Using Itô formula we obtain SDEs for  $e_n$ (We write  $e_n^{\overline{j_1},\overline{j_2},...,\overline{j_m}}(X)$  for an incomplete elementary symmetric polynomial of degree n, written without any of  $X_{j_1},...,X_{j_m}$ .) Methodology of proofs: SDEs for symmetric polynomials of  $dX_i = 2\sqrt{|X_i|}dB_i + \left(\alpha + \sum_{j \neq i} \frac{|X_i| + |X_j|}{X_i - X_i} \mathbb{1}_{\{X_i \neq X_j\}}\right)dt$ 

#### Theorem

If X is a non-colliding solution of BESQ particle system, then  $(e_1, \ldots, e_p)$  are semi-martingales described, for  $n = 1, \ldots, p$ , by

$$de_{n} = \sqrt{\sum_{i=1}^{p} |X_{i}| (e_{n-1}^{\overline{i}})^{2}} dV_{n} + (\sum_{i=1}^{p} \alpha e_{n-1}^{\overline{i}} - \sum_{i < j} (|X_{i}| + |X_{j}|) e_{n-2}^{\overline{i},\overline{j}}) dt$$

Here  $(V_1, \ldots, V_p)$  is a collection of one-dimensional Brownian motions such that

$$d \langle e_n(X), e_m(X) \rangle = 4 \sum_{i=1}^{p} |X_i| e_{n-1}^{\bar{i}}(X) e_{m-1}^{\bar{i}}(X) dt.$$

## Methodology of proofs: SDEs for symmetric polynomials. Non-negative case $0 \le X_1(t) \le \ldots X_p(t)$

#### Theorem

The elementary symmetric polynomials of the non-colliding solution of BESQ particle system starting from  $0 \le x_1 \le ... \le x_p$ are semi-martingales described up to the first exit time  $\tau = \inf\{t > 0 : X_1(t) < 0\}$  by the following system of p SDEs

$$de_n = 2\sqrt{\sum_{k=1}^{p} (2k-1)e_{n-k}e_{n+k-1}}dV_n + (p-n+1)(\alpha-n+1)e_{n-1}dt,$$

where  $V_n$  are one-dimensional Brownian motions such that

$$d\langle e_n, e_m \rangle = 4 \sum_{k=1}^{p} (m-n+2k-1) e_{n-k} e_{m+k-1} dt$$

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McKean argument: non-explosion of  $U = \ln V_N$  since the finite variation part is bounded

EJP(2014): Assumptions (A4-A4') on coefficients of  $d\lambda_i = \sigma_i(\lambda_i)dB_i + \left(b_i(\lambda_i) + \sum_{j \neq i} \frac{H_{ij}(\lambda_i, \lambda_j)}{\lambda_i - \lambda_j}\right)dt$   $i = 1, \dots, p$ 

Conditions for non-collisions

# (A4) In each collision point x there is a force making the particles uncollide

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Conditions for non-collisions

(A4) In each collision point x there is a force making the particles uncollide σ<sup>2</sup><sub>k</sub>(x) + σ<sup>2</sup><sub>l</sub>(x) + H<sub>kl</sub>(x, x) ≠ 0
 Otherwise, if σ<sup>2</sup><sub>k</sub>(x) + σ<sup>2</sup><sub>l</sub>(x) + H<sub>kl</sub>(x, x) = 0, such critical points x are isolated and verify the following condition (A4')

## (A4') In each critical collision point x there is a force making the particles leave from it

In points critical for (A4), the martingale part as well as the force coming from repulsive forces between particles can not move them from x. Then we require

$$\sum_{i=k}^{l} \left( b_i(x) + \sum_{j=1}^{p-2} \frac{H_{ij}(x, y_j)}{x - y_j} \mathbb{1}_{\mathbb{R} \setminus \{x\}}(y_j) \right) \neq 0,$$

for every  $y_1, \ldots, y_{p-2} \in \mathbb{R}$ . It guarantees that the force coming from the whole drift part will move particles from the critical point and cause instant diffraction of the particles.

#### Proof of the Stochastic Gindikin set

Suppose that  $\alpha and <math>\alpha \notin \mathbb{N}$ . Suppose that the particles  $(X_i(t))$  are non-negative.

We compute the expected value of the symmetric polynomials

$$\mathbb{E}e_1(t)=e_1(0)+p\alpha\int_0^t ds=e_1(0)+p\alpha t.$$

$$\begin{split} \mathbb{E}e_2(t) &= e_2(0) + (p-1)(\alpha-1) \int_0^t \mathbb{E}e_1(s) ds \\ &= e_2(0) + (p-1)(\alpha-1)e_1(0)t + p(p-1)\alpha(\alpha-1)\frac{t^2}{2}, \end{split}$$

and so on. Consequently, the coefficient of  $t^n$  in  $\mathbb{E}e_n(t)$  is

$$\frac{1}{n!}p(p-1)\cdot\ldots\cdot(p-n+1)\cdot\alpha(\alpha-1)\cdot\ldots\cdot(\alpha-n+1).$$

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Take *n* the first integer greater than  $\alpha + 1$ . Then  $\mathbb{E}e_n(t)$  is a polynomial with the leading coefficient negative  $\Rightarrow \mathbb{E}e_n(t)$  cannot stay positive for every t > 0, contradiction!

Proof. Let a random variable Y on  $\overline{S}_{p}^{+}$  exist with law  $\Gamma_{p,\Sigma}(\beta,\omega)$ , i.e.  $\mathcal{L}(Y)(u) = (\det(I + \Sigma u))^{-\beta} e^{-\operatorname{Tr}(u(I + \Sigma u)^{-1}\omega)}, \quad u \in \overline{S}_{p}^{+}.$  Proof. Let a random variable Y on  $\bar{S}_{p}^{+}$  exist with law  $\Gamma_{p,\Sigma}(\beta,\omega)$ , i.e.  $\mathcal{L}(Y)(u) = (\det(I + \Sigma u))^{-\beta} e^{-\operatorname{Tr}(u(I + \Sigma u)^{-1}\omega)}, \quad u \in \bar{S}_{p}^{+}.$ Step 1.  $\Rightarrow$  All the laws  $\Gamma_{p,tl}(\beta,\omega'), t \geq 0, rank(\omega') \leq rank(\omega)$  exist

Piotr Graczyk New results on BESQ particle systems and Wishart processes

Proof. Let a random variable Y on  $\overline{S}_p^+$  exist with law  $\Gamma_{p,\Sigma}(\beta,\omega)$ , i.e.  $\mathcal{L}(Y)(u) = (\det(I + \Sigma u))^{-\beta} e^{-\operatorname{Tr}(u(I + \Sigma u)^{-1}\omega)}, \quad u \in \overline{S}_p^+.$ Step 1.  $\Rightarrow All the laws <math>\Gamma_{-1}(\beta,\omega'), t \geq 0, \operatorname{rank}(\omega') \leq \operatorname{rank}(\omega)$  exist

⇒ All the laws  $\Gamma_{p,tl}(\beta, \omega')$ ,  $t \ge 0$ ,  $rank(\omega') \le rank(\omega)$  exist (take exponential family of *Y*, use Fourier-Laplace transform and Lévy continuity theorem)

⇒ there exists an affine Markov process  $Y_t$  with state space  $\bar{S}_p^+ \cap \{M : \operatorname{rank}(M) \leq \max\{\operatorname{rank}(\omega), 2\beta\}\}$  and with law of  $Y_t$  equal to  $\Gamma_{p,2tl}(\beta, \omega)$ . ⇒ there exists an affine Markov process  $Y_t$  with state space  $\bar{S}_p^+ \cap \{M : \operatorname{rank}(M) \leq \max\{\operatorname{rank}(\omega), 2\beta\}\}$  and with law of  $Y_t$  equal to  $\Gamma_{p,2tl}(\beta, \omega)$ .

**Step 2.** The affine Markov process  $Y_t$  is a weak solution of the BESQ Matrix SDE, with  $\alpha = 2\beta$ .